
Decay of viscoelastic waves with memory

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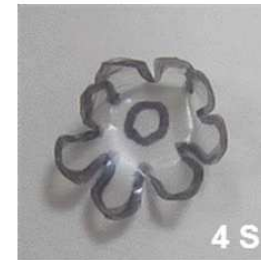
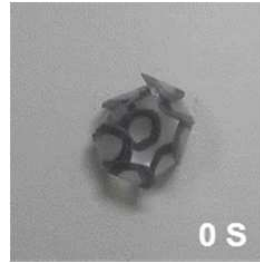
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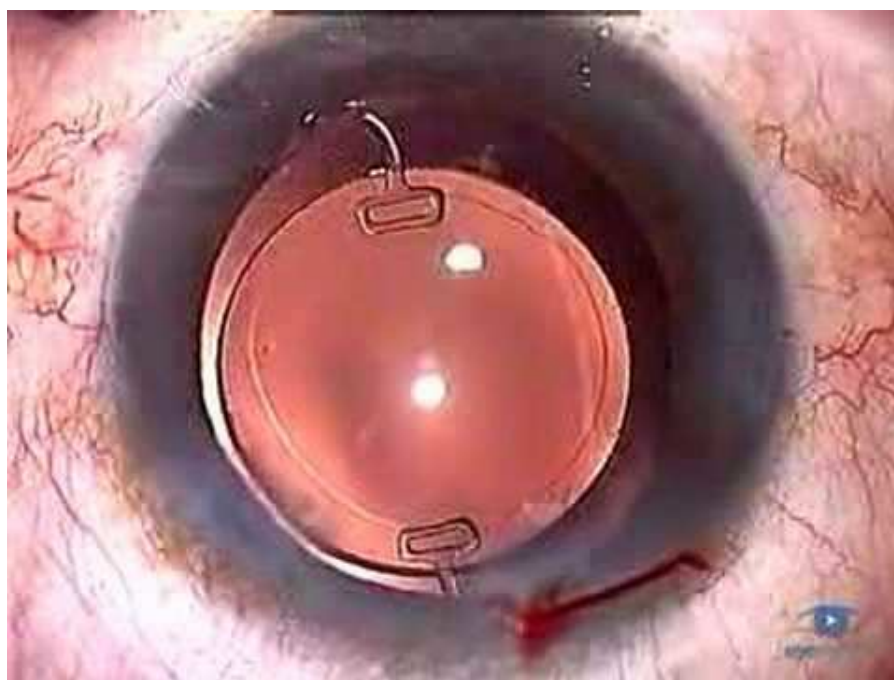
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Summary

- Modelling the displacement
- The energy behavior of the wave equation with memory
- Numerical waves with memory
- Conclusions

Modelling the displacement

- Newton's second law - u displacement of the solid; force; stress

$$\rho \frac{d^2 u}{dt^2} = \nabla \cdot \boldsymbol{\sigma} + f$$

- Stress and strain: Boltzmann representation

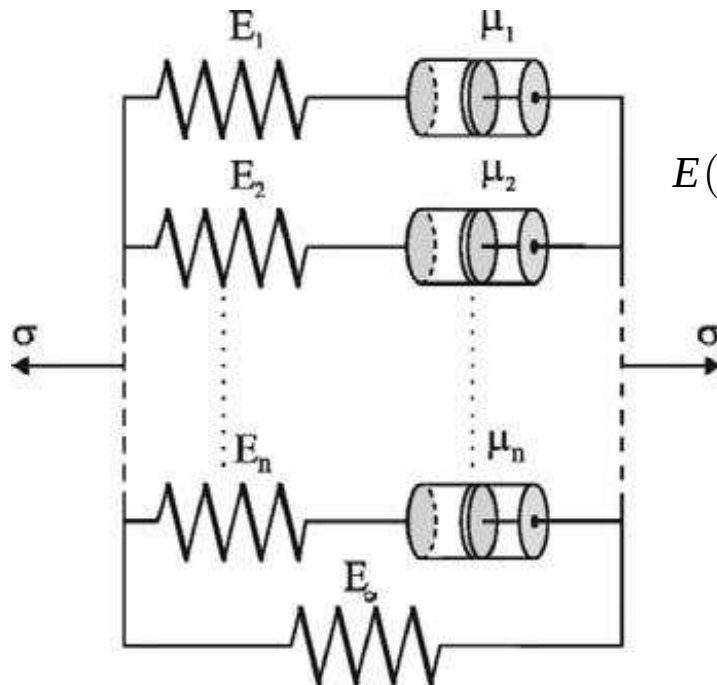
$$\boldsymbol{\sigma}(t) = \underbrace{E(0)\mathbf{D}\boldsymbol{\varepsilon}(t)}_{\text{Hooke's law}} - \underbrace{\int_0^t \frac{\partial}{\partial s} E(t-s)\mathbf{D}\boldsymbol{\varepsilon}(s) ds}_{\text{hereditary integral}}$$

- Strain and displacement:

$$\boldsymbol{\varepsilon}(t) = \frac{1}{2} \left(\nabla u(t) + \nabla u(t)^t \right)$$

$$\rho \frac{d^2 u}{dt^2} - \nabla \cdot \left(\frac{1}{2} E(0) \mathbf{D} (\nabla u(t) + \nabla u(t)^t) \right) = - \int_0^t \frac{\partial}{\partial s} E(t-s) \nabla \cdot \left(\mathbf{D} (\nabla u(s) + \nabla u(s)^t) \right) ds + f$$

Maxwell-Wiechert model:



$$E(t) = E_0 + \sum_{i=1}^n E_i e^{-\frac{E_i}{\mu_i} t}$$

$$\rho \frac{d^2 u}{dt^2} - \hat{E} \nabla \cdot \left(\frac{1}{2} \mathbf{D}(\nabla u(t) + \nabla u(t)^t) \right) = - \int_0^t K_{er}(t-s) \nabla \cdot \left(\mathbf{D}(\nabla u(s) + \nabla u(s)^t) \right) ds + f$$

$$\hat{E} = \sum_{j=0}^n E_j \quad K_{er}(s) = \sum_{i=1}^n \frac{E_i^2}{\mu_i} e^{-\frac{E_i}{\mu_i} s}, \quad u(t) = 0 \text{ on } \partial\Omega$$

General wave equations with memory

$$\frac{d^2u}{dt^2}(t) + c \frac{du}{dt}(t) + \mathcal{A}u(t) = \int_0^t K_{er}(t-s) \mathcal{B}u(s) ds$$

$$\mathcal{A}v = - \sum_{i,j=1}^n \frac{\partial}{\partial x_i} \left(a_{ij} \frac{\partial v}{\partial x_j} \right) + \sum_{i=1}^n \frac{\partial}{\partial x_i} (a_i v) + a_0 v,$$

$$\mathcal{B}v = - \sum_{i,j=1}^n \frac{\partial}{\partial x_i} \left(b_{ij} \frac{\partial v}{\partial x_j} \right) + \sum_{i=1}^n \frac{\partial}{\partial x_i} (b_i v) + b_0 v,$$

Let $u \in H^1(\mathbb{R}^+, H_0^1(\Omega))$ be such that $\frac{d^2u}{dt^2} \in L^2(\mathbb{R}^+, L^2(\Omega))$ and, for all $T > 0$, holds the following

$$\left\{ \begin{array}{l} \left(\frac{d^2u}{dt^2} + c \frac{du}{dt}, w \right) + a(u, w) = \int_0^t K_{er}(t-s) b(u(s), w) ds \text{ in } (0, T), \forall w \in H_0^1(\Omega), \\ \frac{du}{dt}(0) = u_1, \\ u(0) = u_0, \end{array} \right.$$

$$u_0 \in H_0^1(\Omega), u_1 \in L^2\Omega)$$

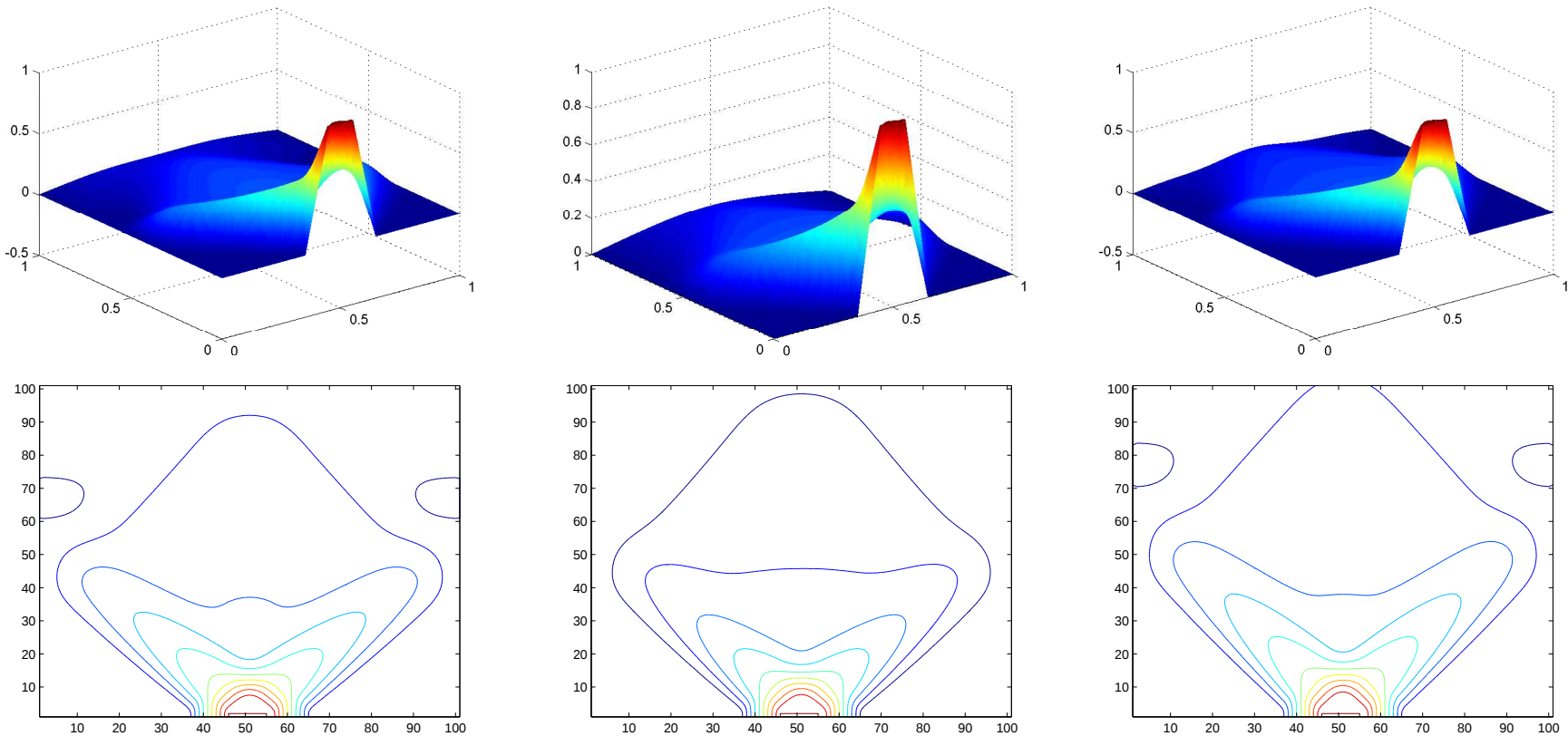
Wave equation with memory- relaxation of a wave equation

Wave equation with memory (WEM): $\rho \frac{d^2u}{dt^2}(t) + c \frac{du}{dt}(t) - D_1 \Delta u(t) = - \int_0^t K_{er}(t-s) D_2 \Delta u(s) ds$

$\tau = \frac{\mu_1}{E_1}, K_{er}(s) = \frac{1}{\tau} e^{-\frac{s}{\tau}}$

■ $\tau \rightarrow 0$

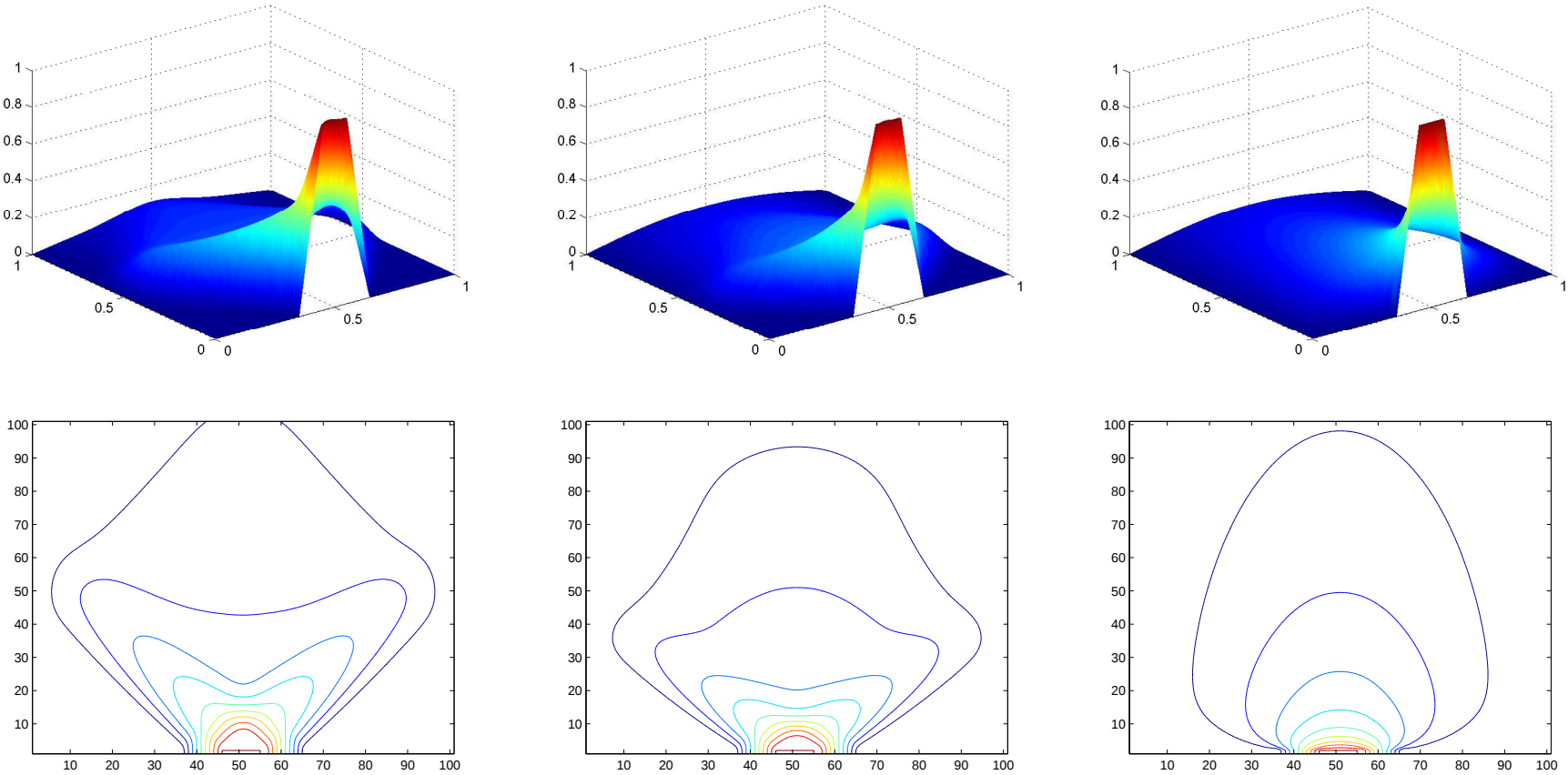
$\rho \frac{d^2u}{dt^2}(t) + c \frac{du}{dt}(t) = (D_1 - D_2) \Delta u(t)$



Wave equation with memory- singular perturbation of a diffusion equation with memory

$\rho \rightarrow 0$

$$\frac{du}{dt}(t) - D_1 \Delta u(t) = - \int_0^t K_{er}(t-s) D_2 \Delta u(s) ds$$



Wave equation

$$\text{Energy} \rightarrow E(t) = \frac{\rho}{2} \left\| \frac{du}{dt}(t) \right\|^2 + \frac{D}{2} \|\nabla u(t)\|^2, t \geq 0$$

- Without damping effect ($c = 0$) - Energy conservation principle

$$E(t) = E(0), t \geq 0$$

- With damping effect ($c > 0$) - $E(t)$ is a decreasing function and

$$E(t) \rightarrow 0, t \rightarrow \infty, \text{exponentially}$$

Diffusion equation with memory

$$\text{Energy} \rightarrow E(t) = \|u(t)\|^2 + \int_0^t \|\nabla u(s)\|^2 ds, t \geq 0$$

- $E(t)$ is bounded in bounded time intervals
- There exists $\gamma > 0$ such that

$$\|u(t)\|^2 + \int_0^t e^{-2\gamma(t-s)} \|\nabla u(s)\|^2 ds \rightarrow 0, t \rightarrow \infty, \text{ exponentially}$$

Wave equation with memory - damping effect

$$\left(\left(\int_0^t K_{er}(t-s) \nabla u(s), \nabla \frac{du}{dt}(t) \right) ds \right.$$

Using the representation

$$\begin{aligned} \left(\int_0^t K_{er}(t-s) \nabla u(s), \nabla \frac{du}{dt}(t) \right) ds &= -\frac{1}{2} \frac{d}{dt} \left\| \int_0^t K_{er}(t-s) \nabla u(s) ds - \nabla u(t) \right\|^2 \\ &\quad + \left\| \int_0^t K'_{er}(t-s) \nabla u(s) ds - \nabla u(t) \right\|^2 + \frac{1}{2} \frac{d}{dt} \left\| \nabla u(t) \right\|^2 \end{aligned}$$

then with $K_{er}(s) = \frac{1}{\tau} e^{-\frac{s}{\tau}}$, we have

$$E(t) + \left\| \int_0^t K_{er}(t-s) \nabla u(s) ds - \nabla u(t) \right\|^2$$

is a decreasing function in time, and consequently the memory term has a damping effect.

Moreover

$$\left\| \int_0^t K_{er}(t-s) \nabla u(s) ds - \nabla u(t) \right\|^2$$

decreases exponentially in time.

The wave equation with memory has an exponential decreasing energy?

Using the representation

$$\begin{aligned} \left(\left(\int_0^t K_{er}(t-s) \nabla u(s), \nabla \frac{du}{dt}(t) \right) ds = -\frac{1}{2} \frac{d}{dt} \int_0^t K_{er}(t-s) \|\nabla u(s) - \nabla u(t)\|^2 ds \right. \\ \left. + \frac{1}{2} \int_0^t K'_{er}(t-s) \|\nabla u(s) - \nabla u(t)\|^2 ds + \frac{1}{2} \int_0^t K_{er}(t-s) \frac{d}{dt} \|\nabla u(t)\|^2 \right. \end{aligned}$$

then with $K_{er}(s) = \frac{1}{\tau} e^{-\frac{s}{\tau}}$, we have

$$E(t) - \int_0^t K_{er}(t-s) ds \|\nabla u(t)\|^2 + \int_0^t K_{er}(t-s) \|\nabla u(s) - \nabla u(t)\|^2 ds$$

is a decreasing function in time, and

$$\|\nabla u(t)\|^2 + \int_0^t K_{er}(t-s) \|\nabla u(s) - \nabla u(t)\|^2 ds$$

decreases exponentially in time.

The kinetic energy of the wave equation with memory decreases exponentially in time?

Kinetic energy $\| \frac{du}{dt}(t) \|^2$

Wave equation with damping effect:

$$\frac{d^2 u}{dt^2} + c \frac{du}{dt} + \beta u(t) = D \Delta u(t)$$

Then

$$E(t) + \|u(s)\|^2 \rightarrow 0, t \rightarrow \infty, \text{ exponentially.}$$

Wave equation with memory and damping effect:

$$\frac{d^2 u}{dt^2} + c \frac{du}{dt} + \beta u(t) - D_1 \Delta u(t) = - \int_0^t K_{er}(t-s) D_2 \Delta u(s) ds$$

Then, for $K_{er}(s) = \frac{1}{\tau} e^{-\frac{s}{\tau}}$, $\exists \gamma > 0$ such that

$$E(t) + \|u(t)\|^2 + \left\| \int_0^t K_{er}(t-s) e^{-\gamma(t-s)} \nabla u(s) ds - \nabla u(t) \right\|^2 \rightarrow 0, t \rightarrow \infty, \text{ exponentially.}$$

General wave equations with memory

$$\begin{aligned} \int_0^t K_{er}(t-s)b(u(s), \frac{du}{dt}(t)) ds &= \frac{d}{dt} \int_0^t K_{er}(t-s)b(u(s), u(t)) ds \\ &\quad - K_{er}(0)b(u(t), u(t)) - \int_0^t K'_{er}(t-s)b(u(s), u(t)) ds \\ &\quad - \int_0^t K_{er}(t-s)b(u(s), \frac{du}{dt}(t)) ds \end{aligned}$$

$a(.,.), b(.,.)$ are H^1 – elliptic, $b(.,.)$ is bounded, $K_{er} \in H^1(\mathbb{R}^+)$

- If $c = 0$, then $E(t)$ is bounded in bounded time intervals
- If $c, \beta > 0$, then there exists a family of wave problems such that

$$E(t) + \|u(t)\|^2 \rightarrow 0, t \rightarrow \infty, \text{ exponentially.}$$

Numerical waves and their qualitative behavior

- $\mathcal{T}_h$ triangulation of $\Omega \subset \mathbb{R}^2$ with diameter h
- finite element space

$$\mathcal{V}_h = \{v \in C^0(\overline{\Omega}) : v = 0 \text{ on } \partial\Omega, v \in \mathbb{P}(K), K \in \mathcal{T}_h\}$$

$$\left\{ \begin{array}{l} \left(\frac{d^2 u_h}{dt^2} + c \frac{du_h}{dt}, w_h \right) + a(u_h, w_h) = \int_0^t K_{er}(t-s) b(u_h(s), w_h) ds \text{ in } (0, T), \forall w_h \in \mathcal{V}_h, \\ \frac{du_h}{dt}(0) = u_{1,h}, \\ u_h(0) = u_{0,h}. \end{array} \right.$$

Energy estimates: numerical waves mimic their continuous counterpart

$$\mathcal{A}u = -\nabla \cdot (D_1 \nabla u), \mathcal{B}u = -\nabla \cdot (D_2 \nabla u)$$

- Discretize in space (FD+trapz)

$$\begin{aligned} & \left(\frac{u_h^{n+1} - 2u_h^n + u_h^{n-1}}{\Delta t^2}, v \right) + \left(c \frac{u_h^{n+1} - u_h^{n-1}}{2\Delta t}, v \right) + (D_1 \nabla u_h^{n+1}, \nabla v) \\ &= \underbrace{\frac{\Delta t}{2} \sum_{j=0}^n (K_{er}(t_{n+1} - t_{j+1}) D_2 \nabla u_h^{j+1} + K_{er}(t_{n+1} - t_j) D_2 \nabla u_h^j, \nabla v)}_{I_{n+1}} \end{aligned}$$

- $K_{er}(s) = \frac{1}{\tau} e^{-\frac{s}{\tau}}$

$$\begin{aligned} & \left(\left(\frac{1}{\Delta t^2} + \frac{c}{2\Delta t} \right) u_h^{n+1}, v \right) + \left(\left(D_1 - \frac{\Delta t}{2\tau} D_2 \right) \nabla u_h^{n+1}, \nabla v \right) \\ &= \left(\left(\frac{2}{\Delta t^2} + \frac{\Delta t}{2\tau} e^{-\frac{\Delta t}{\tau}} D_2 \nabla \right) u_h^n, v \right) + \left(\left(\frac{c}{2\Delta t} - \frac{1}{\Delta t^2} \right) u_h^{n-1}, v \right) + e^{-\frac{\Delta t}{\tau}} (I_n, \nabla v) \end{aligned}$$

- Known data: u_h^0, u_h^1 and I_1

- Given u_h^n and I_n , do

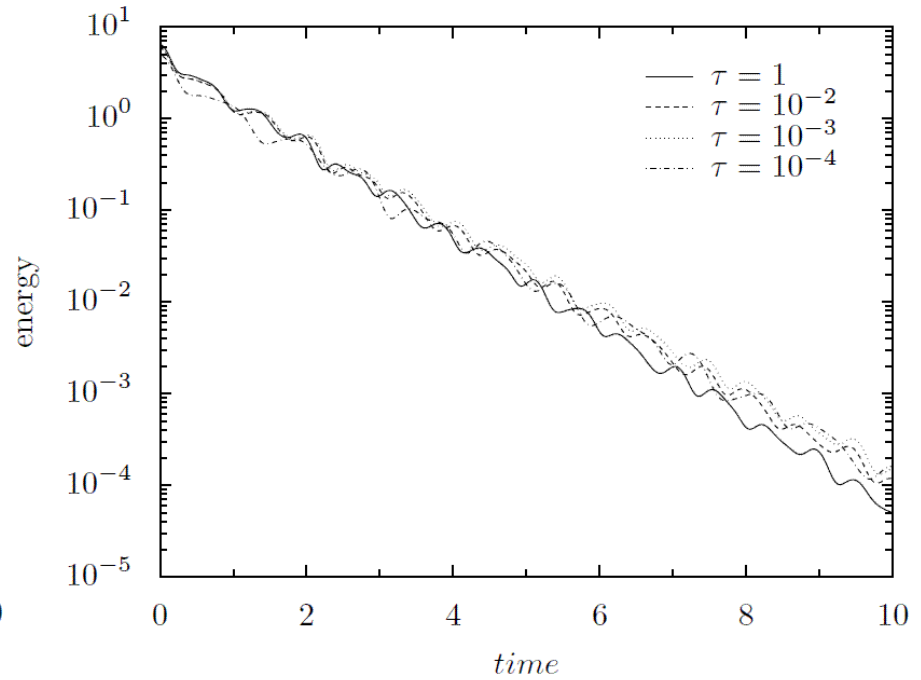
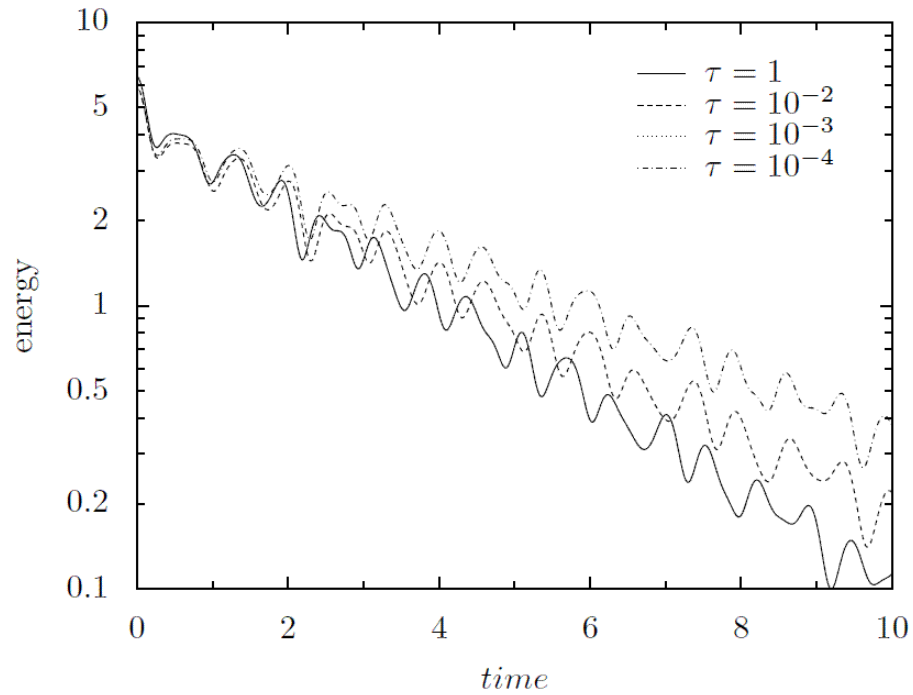
- Solve for u_h^{n+1}

- Update $I_{n+1} : I_{n+1} = e^{-\frac{\Delta t}{\tau}} I_n + \frac{D_2 \Delta t}{2\tau} (e^{-\frac{\Delta t}{\tau}} \nabla u_h^n + \nabla u_h^{n+1})$

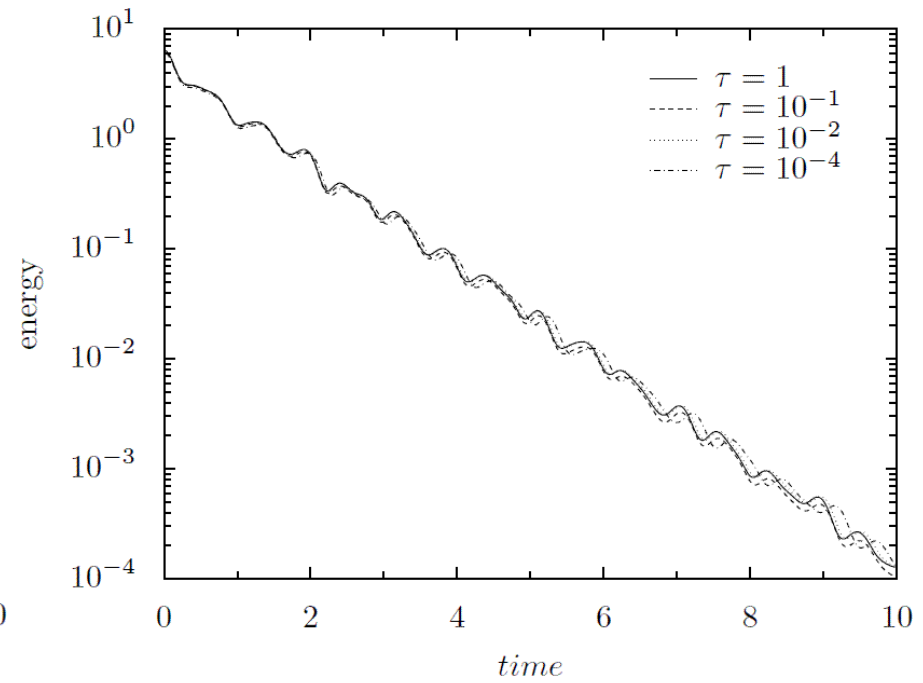
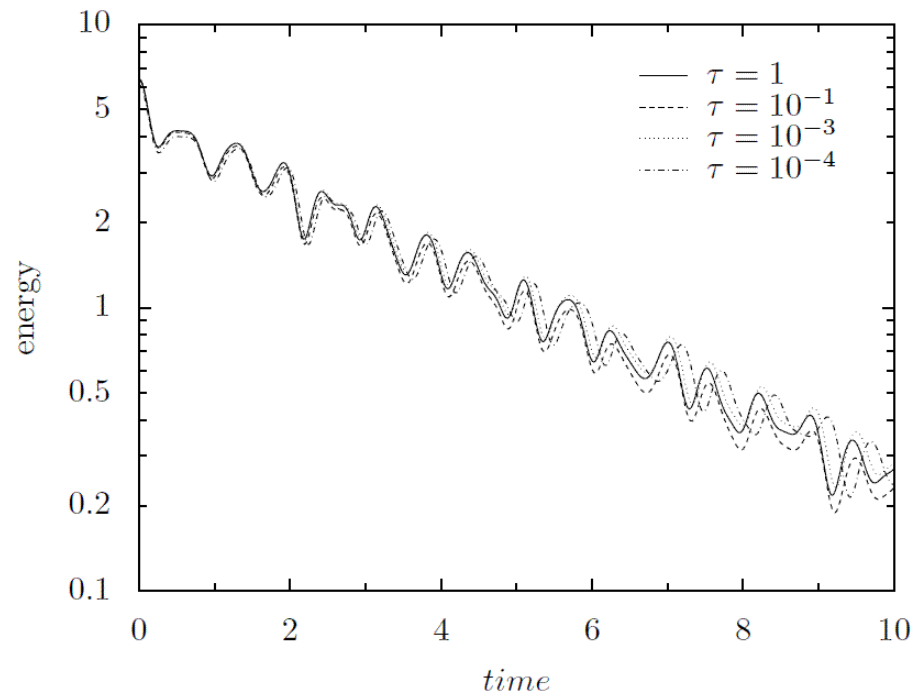
Behavior of the (discrete) energy

$$E_{h,n} = \left\| \frac{u_h^n - u_h^{n-2}}{2\Delta t} \right\|^2 + \|u_h^n\|_1^2 + \|I_n - \nabla u_h^n\|^2, \quad n \geq 2.$$

$$\Omega = (-1, 1)^2, u_0(x, y) = e^{-\frac{x^2+y^2}{0.1}}$$



Effect of memory coefficient τ for different damping effects ($D_2 = 0.1, c = 0.25, 1$)



Effect of memory coefficient τ for different damping effects ($D_2 = 0.01, c = 0.25, 1$)

The discrete energy decays exponentially to zero.

Conclusions:

- The energy of the wave equation satisfies the energy conservation principle. If a damping effect is presented ($c > 0$) then $E(t)$ decreases exponentially in time.
- The memory term in the wave equation with memory (WEM) ($c = 0$) has a damping effect - the potential and kinetic energies decrease in time and the potential energy decreases exponentially in time.
- The energy of the WEM with $c, \beta > 0$, decreases exponentially in time.
- The last result can be generalized for a general class of wave equations with memory.