

Exercise 1 – p -brane action

Consider a p -brane moving in a spacetime (M, g) , which we take to be d -dimensional Minkowski spacetime $(\mathbb{R}^{1,d-1}, \eta)$. The action for this p -brane is given by

$$S = -\frac{T_p}{2} \int_{\Sigma} d^{p+1} \sigma \sqrt{-h} h^{\alpha\beta} \eta_{\mu\nu} \partial_{\alpha} X^{\mu} \partial_{\beta} X^{\nu} + \Lambda_p \int_{\Sigma} d^{p+1} \sigma \sqrt{-h} .$$

Here, the embedding $X : (\Sigma, h_{\alpha\beta}) \rightarrow (M, \eta_{\mu\nu})$ denotes a smooth map between two semi-Riemannian manifolds, the $(p+1)$ -dimensional world volume of the p -brane and d -dimensional Minkowski spacetime, and $h \equiv \det h_{\alpha\beta}$.

Investigate whether the equation of motion for the world-volume metric $h_{\alpha\beta}$ is solved by equating the world-volume metric with the induced metric $X^* \eta$. Show that this requires a non-vanishing cosmological constant Λ_p for $p \neq 1$.

Exercise 2 – Weyl rescalings in two dimensions

Consider a two-dimensional semi-Riemannian manifold $(M, h_{\alpha\beta})$.

a) Show that under local Weyl rescalings of the metric, $h_{\alpha\beta} \rightarrow e^{2\Lambda(\sigma)} h_{\alpha\beta}$, the quantity $\sqrt{-h} \mathcal{R}$ transforms as

$$\sqrt{-h} \mathcal{R} \rightarrow \sqrt{-h} (\mathcal{R} - 2\nabla^2 \Lambda) . \quad (1)$$

Here, $h \equiv \det h_{\alpha\beta}$, and \mathcal{R} denotes the Ricci scalar ($\mathcal{R} = h^{\alpha\beta} \text{Ric}_{\alpha\beta}$, $\text{Ric}_{\alpha\beta} = R^{\gamma}_{\alpha\gamma\beta}$).

b) Conclude that one may add a term $\int_{\Sigma} d^2 \sigma \sqrt{-h} \mathcal{R}$ to the world-sheet action of a closed string while maintaining invariance under local Weyl rescalings.

c) Use the result (1) to show that locally, every metric of signature $(-1, 1)$ can be brought into the form $\eta = \text{diag}(-1, 1)$ by Weyl rescalings and appropriate choice of local coordinates.

Exercise 3 – Witt algebra

In the conformal gauge, and in light-cone coordinates $\sigma_{\pm} = \tau \pm \sigma$, the world-sheet energy-momentum tensor $T_{\alpha\beta}$ has components given by $T_{+-} = T_{-+} = 0$, $T_{\pm\pm} = -\partial_{\pm} X^{\mu} \partial_{\pm} X^{\nu} \eta_{\mu\nu}$, where $\partial_{\pm} = \partial_{\sigma\pm}$.

Consider the closed string.

a) Using the equal time Poisson brackets $\{\cdot, \cdot\}$,

$$\{X^{\mu}(\sigma, \tau), X^{\nu}(\sigma', \tau)\} = \{\dot{X}^{\mu}(\sigma, \tau), \dot{X}^{\nu}(\sigma', \tau)\} = 0, \quad \{X^{\mu}(\sigma, \tau), \dot{X}^{\nu}(\sigma', \tau)\} = 2\pi\alpha' \eta^{\mu\nu} \delta(\sigma - \sigma'),$$

calculate

$$\{T_{\pm\pm}(\sigma, \tau), X^{\mu}(\sigma', \tau)\} .$$

b) On-shell, $T_{--} = T_{--}(\sigma^-)$, $T_{++} = T_{++}(\sigma^+)$.

Define charges

$$L_{\epsilon^-} \equiv -\frac{1}{2\pi} \int_0^{2\pi} d\sigma \epsilon^-(\sigma^-) T_{--}(\sigma^-) ,$$

with $\epsilon^-(\sigma^-)$ a periodic function in σ . Show that L_{ϵ^-} is conserved and

$$\{L_{\epsilon^-}, X^\mu(\sigma, \tau)\} = -\alpha' \mathcal{L}_{\epsilon^-} X^\mu(\sigma, \tau) ,$$

where \mathcal{L} denotes the Lie derivative.

Using this result, and Fourier decomposing

$$T_{--}(\sigma^-) = -\alpha' \sum_{n \in \mathbb{Z}} L_n e^{-im\sigma^-} ,$$

show that

$$\{L_m, X^\mu(\sigma, \tau)\} = T_m X^\mu(\sigma, \tau) ,$$

where the vector field T_m is an element of the Lie algebra

$$[T_m, T_n] = i(m - n) T_{m+n} .$$

c) Show that the generators L_m satisfy the Witt algebra

$$\{L_m, L_n\} = -i(m - n) L_{m+n}$$

Verify that these commutation relations satisfy the Jacobi identity.