Riemannian Geometry Exam - January 30, 2018 LMAC and MMA

Solutions

1.

a) G is a subgroup of $\mathcal{M}_{2\times 2}(\mathbb{C})$. Indeed, suppose $A \in G$ and $B \in G$. Then $AB \in G$ since

$$(AB)^*J(AB) = B^*(A^*JA)B = B^*JB = J,$$

and $A^{-1} \in G$ as

$$A^*JA = J \Rightarrow (A^*)^{-1}A^*JAA^{-1} = (A^{-1})^*JA^{-1} \Rightarrow J = (A^{-1})^*JA^{-1}.$$

b) We denote by $S_{2\times 2}(\mathbb{C})$ the space of skew-hermitean 2×2 matrices. This is a 4-dimensional space. Let $f: \mathcal{M}_{2\times 2} \to \mathcal{S}_{2\times 2}$ be defined by

$$f(A) = A^* J A.$$

This function is smooth and

$$Df(A)(B) = B^*JA + A^*JB.$$

Suppose $A \in f^{-1}(J)$ and $C \in T_J \mathcal{S}_{2 \times 2}(\mathbb{C}) \equiv \mathcal{S}_{2 \times 2}(\mathbb{C})$. Choosing $B = -\frac{1}{2}AJC$, we get

$$Df(A)\left(-\frac{1}{2}AJC\right) = \frac{1}{2}C^*JA^*JA - \frac{1}{2}A^*JAJC$$
$$= -\frac{1}{2}CJ^2 - \frac{1}{2}J^2C = C.$$

This shows that f is a submersion at A. Since A is arbitrary in $f^{-1}(J)$, J is a regular value of f. It follows that $G = f^{-1}(J)$ is a submanifold of $\mathcal{M}_{2\times 2}(\mathbb{C})$ of dimension 8 - 4 = 4. Hence, G is a Lie group. c)

$$T_I G = \ker Df(I) = \{ B \in \mathcal{M}_{2 \times 2} : B^*J + JB = 0 \}.$$

A basis for $T_I G$ is $\{B_1, B_2, B, \tilde{B}\}$, where

$$B_1 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}, \quad B_2 = \begin{bmatrix} i & 0 \\ 0 & i \end{bmatrix}, \quad B = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}, \quad \tilde{B} = \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}.$$

One easily checks that $[B_1, B] = 2B$, $[B_1, \tilde{B}] = 2\tilde{B}$ and $[B, \tilde{B}] = B_1$. Obviously, the bracket of B_2 with any other matrix is 0. d) X and \tilde{X} are the left invariant vector fields that take the values B and \tilde{B} at the identity, respectively. Therefore, $[X, \tilde{X}]$ is the left invariant vector field that takes the value equal to the bracket of B and \tilde{B} at the identity, i.e.

$$[X, \tilde{X}]_A = A[B, \tilde{B}] = AB_1.$$

2.

a) $\partial_x = \left(1, 0, -\frac{y}{x^2+y^2}\right)$ and $\partial_y = \left(0, 1, \frac{x}{x^2+y^2}\right)$. The unit normal to \mathcal{H} with positive third component is

$$n = \frac{\partial_x \times \partial_y}{\|\partial_x \times \partial_y\|} = \frac{\left(\frac{y}{\sqrt{x^2 + y^2}}, -\frac{x}{\sqrt{x^2 + y^2}}, \sqrt{x^2 + y^2}\right)}{\sqrt{x^2 + y^2 + 1}}$$

b)

$$\begin{split} \iota(n)(dx \wedge dy \wedge dz) &= (\iota(n)dx) \wedge dy \wedge dz - dx \wedge (\iota(n)dy) \wedge dz \\ &+ dx \wedge dy \wedge (\iota(n)dz) \\ &= \frac{y}{\sqrt{x^2 + y^2}\sqrt{x^2 + y^2 + 1}} \, dy \wedge dz \\ &+ \frac{x}{\sqrt{x^2 + y^2}\sqrt{x^2 + y^2 + 1}} \, dx \wedge dz \\ &+ \frac{\sqrt{x^2 + y^2}}{\sqrt{x^2 + y^2 + 1}} \, dx \wedge dy = \omega. \end{split}$$

c)

$$\eta = p^* \omega = \frac{1}{\sqrt{r^2 + 1}} \left(\sin \theta \left| \begin{array}{c} \sin \theta & r \cos \theta \\ 0 & 1 \end{array} \right| + \\ + \cos \theta \left| \begin{array}{c} \cos \theta & -r \sin \theta \\ 0 & 1 \end{array} \right| + \\ + r \left| \begin{array}{c} \cos \theta & -r \sin \theta \\ \sin \theta & r \cos \theta \end{array} \right| \right) dr \wedge d\theta \\ = \sqrt{r^2 + 1} dr \wedge d\theta.$$

d)

$$L_X \eta = L_{\left(\frac{\partial f}{\partial r}\partial_r + \frac{1}{r^2 + 1}\frac{\partial f}{\partial \theta}\partial_\theta\right)} \sqrt{r^2 + 1} \, dr \wedge d\theta$$

$$= \frac{\partial f}{\partial r} \frac{r}{\sqrt{r^2 + 1}} \, dr \wedge d\theta$$

$$+ \sqrt{r^2 + 1} \frac{\partial^2 f}{\partial r^2} \, dr \wedge d\theta$$

$$+ \frac{1}{\sqrt{r^2 + 1}} \frac{\partial^2 f}{\partial \theta^2} \, dr \wedge d\theta$$

$$= (\operatorname{div} X) \eta,$$

where

div
$$X = \frac{\partial f}{\partial r} \frac{r}{r^2 + 1} + \frac{\partial^2 f}{\partial r^2} + \frac{1}{r^2 + 1} \frac{\partial^2 f}{\partial \theta^2}$$

e) Using Cartan's formula and Stokes' Theorem, we have

$$\int_{S} L_{X} \eta = \int_{S} [d(\iota(X)\eta) + \iota(X)d\eta] = \int_{\partial S} \iota(X)\eta$$

$$= \int_{\partial S} \iota\left(\frac{\partial f}{\partial r}\partial_{r} + \frac{1}{r^{2}+1}\frac{\partial f}{\partial \theta}\partial_{\theta}\right)\sqrt{r^{2}+1}\,dr \wedge d\theta$$

$$= \int_{\partial S} -\frac{1}{\sqrt{r^{2}+1}}\frac{\partial f}{\partial \theta}\,dr + \frac{\partial f}{\partial r}\sqrt{r^{2}+1}\,d\theta$$

$$= \int_{\theta_{0}}^{\theta_{1}}\frac{\partial f}{\partial r}(r_{1},\theta)\sqrt{r_{1}^{2}+1}\,d\theta - \int_{\theta_{0}}^{\theta_{1}}\frac{\partial f}{\partial r}(r_{0},\theta)\sqrt{r_{0}^{2}+1}\,d\theta$$

$$+ \int_{r_{0}}^{r_{1}}\frac{1}{\sqrt{r^{2}+1}}\frac{\partial f}{\partial \theta}(r,\theta_{1})\,dr - \int_{r_{0}}^{r_{1}}\frac{1}{\sqrt{r^{2}+1}}\frac{\partial f}{\partial \theta}(r,\theta_{0})\,dr.$$

3.

a) Since $(x, y, z) = (r \cos \theta, r \sin \theta, \theta)$, clearly

$$\begin{cases} dx = \cos\theta \, dr - r\sin\theta \, d\theta, \\ dy = \sin\theta \, dr + r\cos\theta \, d\theta, \\ dz = d\theta. \end{cases}$$

So the metric induced on the helicoid by the euclidean metric,

$$ds^2 = dx^2 + dy^2 + dz^2,$$

is

$$ds^2 = dr^2 + (r^2 + 1) \, d\theta^2.$$

b) The lagrangian is

$$L(r,\theta,\dot{r},\dot{\theta}) = \frac{1}{2} \left(\dot{r}^2 + (r^2 + 1)\dot{\theta}^2 \right).$$

The equations for the geodesics are

$$\frac{d}{dt}\frac{\partial L}{\partial \dot{r}} - \frac{\partial L}{\partial r} = 0,
\frac{d}{dt}\frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = 0.$$

These equations are equivalent to

$$\begin{aligned} \ddot{r} - r\theta^2 &= 0, \\ \frac{d}{dt}[(r^2 + 1)\dot{\theta}] &= 0 \iff \ddot{\theta} + \frac{2r}{r^2 + 1}\dot{r}\dot{\theta} &= 0. \end{aligned}$$

We read that the nonzero Christoffel symbols are

$$\Gamma^r_{\theta\theta} = -r, \qquad \Gamma^{\theta}_{r\theta} = \Gamma^{\theta}_{\theta r} = \frac{r}{r^2 + 1}.$$

c) We know that

$$\nabla_{\partial_r}\partial_r = 0, \quad \nabla_{\partial_r}\partial_\theta = \nabla_{\partial_\theta}\partial_r = \frac{r}{r^2 + 1}\partial_\theta, \quad \nabla_{\partial_\theta}\partial_\theta = -r\partial_r.$$

Using these equalities, we get

$$R(\partial_r, \partial_\theta)\partial_r = \nabla_{\partial_r} \nabla_{\partial_\theta} \partial_r - \nabla_{\partial_\theta} \nabla_{\partial_r} \partial_r - \nabla_{[\partial_r, \partial_\theta]} \partial_r$$

$$= \nabla_{\partial_r} \left(\frac{r}{r^2 + 1} \partial_\theta \right)$$

$$= \frac{1 - r^2}{(r^2 + 1)^2} \partial_\theta + \frac{r^2}{(r^2 + 1)^2} \partial_\theta$$

$$= \frac{1}{(r^2 + 1)^2} \partial_\theta.$$

Thus

$$R(\partial_r, \partial_\theta, \partial_r, \partial_\theta) = (R(\partial_r, \partial_\theta)\partial_r, \partial_\theta) = \frac{1}{r^2 + 1}.$$

The curvature of \mathcal{H} is

$$K = -\frac{R(\partial_r, \partial_\theta, \partial_r, \partial_\theta)}{(\partial_r, \partial_r)(\partial_\theta, \partial_\theta) - (\partial_r, \partial_\theta)^2} = -\frac{1}{(r^2 + 1)^2}.$$

d) Since V is parallel along c and $\dot{c} = \partial_r$,

$$\nabla_{\partial_r} (V^r \partial_r + V^\theta \partial_\theta) = \frac{\partial V^r}{\partial r} \partial_r + \frac{\partial V^\theta}{\partial r} \partial_\theta + V^\theta \frac{r}{r^2 + 1} \partial_\theta = 0.$$

The components of V satisfy

$$\begin{cases} \frac{\partial V^r}{\partial r} = 0, \\ \frac{\partial V^{\theta}}{\partial r} + \frac{r}{r^2 + 1} V^{\theta} = 0. \end{cases}$$

The solution of this system is

$$\begin{cases} V^{r}(r,\theta_{0}) = V^{r}(r_{0},\theta_{0}), \\ V^{\theta}(r,\theta_{0}) = \frac{\sqrt{r_{0}^{2}+1}}{\sqrt{r^{2}+1}} V^{\theta}(r_{0},\theta_{0}) \end{cases}$$

So, we have

$$V(r,\theta) = V^r(r_0,\theta_0)\partial_r + \frac{\sqrt{r_0^2 + 1}}{\sqrt{r^2 + 1}} V^\theta(r_0,\theta_0)\partial_\theta.$$

Note that the angle between V and \dot{c} is constant (because c is a geodesic) and that the length of V is constant.

e) The reparemeterization of c by arclength is $\tilde{c}(\theta) = \left(r_0, \frac{\theta}{\sqrt{r_0^2 + 1}}\right)$ since $\|\dot{\tilde{c}}\| = \left\|\frac{\partial_{\theta}}{\sqrt{r_0^2 + 1}}\right\| = 1$. The acceleration of \tilde{c} is $\nabla_{\dot{\tilde{c}}}\dot{\tilde{c}} = \nabla_{\frac{\partial_{\theta}}{\sqrt{r_0^2 + 1}}} \frac{\partial_{\theta}}{\sqrt{r_0^2 + 1}} = -\frac{r_0}{r_0^2 + 1}\partial_r.$

The frame $(\dot{\tilde{c}}, -\partial_r)$ is orthonormal and positively oriented. Hence, the geodesic curvature of \tilde{c} is $\frac{r_0}{r_0^2+1}$.

f) We have

$$\partial_r = (\cos\theta, \sin\theta, 0) = \partial_r,$$
$$\partial_\theta = (-r\sin\theta, r\cos\theta, 1) = \tilde{\partial}_\theta + \tilde{\partial}_z.$$

Since the metric is

$$ds^2 = dr^2 + r^2 d\theta^2 + dz^2.$$

 $\tilde{\partial}_r, \frac{1}{r}\tilde{\partial}_{\theta}$ and $\tilde{\partial}_z$ are unit vectors. We note that $\tilde{\partial}_r \times \frac{1}{r}\tilde{\partial}_{\theta} = \tilde{\partial}_z$ and $\tilde{\partial}_r \times \tilde{\partial}_z = -\frac{1}{r}\tilde{\partial}_{\theta}$. We have

$$E_r = \frac{\partial_r}{\|\partial_r\|} = \tilde{\partial}_r,$$
$$E_\theta = \frac{\partial_\theta}{\|\partial_\theta\|} = \frac{\tilde{\partial}_\theta + \tilde{\partial}_z}{\|\tilde{\partial}_\theta + \tilde{\partial}_z\|} = \frac{\tilde{\partial}_\theta + \tilde{\partial}_z}{\sqrt{r^2 + 1}}.$$

The unit normal n to \mathcal{H} with a positive z component is

$$n = E_r \times E_\theta = \frac{-\tilde{\partial}_\theta + r^2 \tilde{\partial}_z}{r\sqrt{r^2 + 1}}.$$

 $\mathbf{g})$

$$\begin{split} \tilde{\nabla}_{E_r} n &= \tilde{\nabla}_{\tilde{\partial}_r} \left(\frac{-\tilde{\partial}_{\theta} + r^2 \tilde{\partial}_z}{r \sqrt{r^2 + 1}} \right) \\ &= \frac{1}{r^2 \sqrt{r^2 + 1}} \tilde{\partial}_{\theta} + \frac{1}{\sqrt{(r^2 + 1)^3}} \tilde{\partial}_{\theta} - \frac{1}{r^2 \sqrt{r^2 + 1}} \tilde{\partial}_{\theta} \\ &+ \frac{1}{\sqrt{r^2 + 1}} \tilde{\partial}_z - \frac{r^2}{\sqrt{(r^2 + 1)^3}} \tilde{\partial}_z \\ &= \frac{1}{\sqrt{(r^2 + 1)^3}} \tilde{\partial}_{\theta} + \frac{1}{\sqrt{(r^2 + 1)^3}} \tilde{\partial}_z, \\ \tilde{\nabla}_{E_{\theta}} n &= \tilde{\nabla}_{\frac{\tilde{\partial}_{\theta} + \tilde{\partial}_z}{\sqrt{r^2 + 1}}} \left(\frac{-\tilde{\partial}_{\theta} + r^2 \tilde{\partial}_z}{r \sqrt{r^2 + 1}} \right) \\ &= \frac{1}{r^2 + 1} \partial_r. \end{split}$$

h) Using the results obtained above, we get

$$\begin{pmatrix} -\tilde{\nabla}_{E_r}n, E_r \end{pmatrix} = 0, \begin{pmatrix} -\tilde{\nabla}_{E_r}n, E_\theta \end{pmatrix} = -\frac{1}{(r^2+1)^2} (\tilde{\partial}_\theta + \tilde{\partial}_z, \tilde{\partial}_\theta + \tilde{\partial}_z) = -\frac{1}{r^2+1}, \begin{pmatrix} -\tilde{\nabla}_{E_\theta}n, E_r \end{pmatrix} = -\frac{1}{r^2+1}, \begin{pmatrix} -\tilde{\nabla}_{E_\theta}n, E_\theta \end{pmatrix} = 0.$$

The second fundamental form of ${\mathcal H}$ is

$$II(r,\theta) = \begin{bmatrix} 0 & -\frac{1}{r^2+1} \\ -\frac{1}{r^2+1} & 0 \end{bmatrix}.$$

i) The principal directions are the eigenvectors of the second fundamental form and the principal curvatures are the eigenvalues of the second fundamental form. These are $E_r + E_{\theta}$ with principal curvature $-\frac{1}{r^2+1}$ and $E_r - E_{\theta}$ with principal curvature $\frac{1}{r^2+1}$ The mean curvature of \mathcal{H} is the average of the principal curvatures, or half the trace of the second fundamental form. In this case the mean curvature is 0.