On the partial regularity of suitable weak solutions: Pointwise estimates and space-time decay

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The subject of the talk is part of two papers with P. Maremonti:

F. C. and P. Maremonti, A remark on the partial regularity of a suitable weak solution to the Navier-Stokes Cauchy problem, submitted.

F. C. and P. Maremonti, *On the spatial asymptotic decay of a suitable weak solution to the Navier-Stokes Cauchy problem*, submitted.

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Introduction

We study the Navier-Stokes Cauchy problem:

$$\begin{aligned} & v_t + v \cdot \nabla v + \nabla \pi_v = \Delta v, \ \nabla \cdot v = 0, \ \text{in} \ (0, T) \times \mathbb{R}^3, \\ & v(0, x) = v_\circ(x) \text{ on } \{0\} \times \mathbb{R}^3. \end{aligned}$$
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It is well known that in

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Theorem 1 (Here $J^2(\mathbb{R}^3)$:=completion of $\mathscr{C}_0(\mathbb{R}^3)$ in L^2 -norm)

For all $v_{\circ} \in J^{2}(\mathbb{R}^{3})$ there exists a suitable weak solution.

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Definition 1 (Here $J^{1,2}(\mathbb{R}^3)$:=completion of $\mathscr{C}_0(\mathbb{R}^3)$ in $W^{1,2}(\mathbb{R}^3)$)

A pair (v, π_v) is said a suitable weak solution to problem (1) if

i) for all T > 0, $v \in L^2(0, T; J^{1,2}(\mathbb{R}^3))$ and $\pi_v \in L^{\frac{5}{3}}((0, T) \times \mathbb{R}^3)$, $\|v(t)\|_2^2 + 2\int_0^t \|\nabla v(\tau)\|_2^2 d\tau \le \|v_\circ\|_2^2, t > 0$.

ii)
$$\lim_{t\to 0} \|v(t) - v_{\circ}\|_{2} = 0,$$

iii) for all
$$t, s \in (0, T)$$
 and for all $\varphi \in C_0^1([0, T) \times \mathbb{R}^3)$:
$$\int_s^t \left[(v, \varphi_\tau) - (\nabla v, \nabla \varphi) + (v \cdot \nabla \varphi, v) + (\pi_v, \nabla \cdot \varphi) \right] d\tau + (v(s), \varphi(s)) = (v(t), \varphi(t));$$

$$\begin{split} \text{iv)} \quad & \text{for all } t \geq s, \text{ for } s = 0 \text{ and } a.e. \text{ in } s \geq 0, \text{ and for all nonnegative } \phi \in C_0^\infty(\mathbb{R} \times \mathbb{R}^3): \\ & \int_{\mathbb{R}^3} |v(t)|^2 \phi(t) + 2 \int_{s-\mathbb{R}^3}^t \int |\nabla v|^2 \phi \leq \int_{\mathbb{R}^3} |v(s)|^2 \phi(s) dx \\ & + \int_{s-\mathbb{R}^3}^t \int |v|^2 (\phi_\tau + \Delta \phi) dx d\tau + \int_{s-\mathbb{R}^3}^t \int (|v|^2 + 2\pi_v) v \cdot \nabla \phi dx d\tau. \end{split}$$

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Partial regularity of suitable weak solutions

It is well known that the regularity of a suitable weak solution and its uniqueness are open problems.

However in [CKN82] and, subsequently in several papers,

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Following [CKN82] we set

$$M(r) := r^{-2} \iint_{Q_r} (|v|^3 + |v||\pi_v|) dy d\tau + r^{-\frac{13}{4}} \int_{t-r^2}^t (\int_{|x-y| < r} |\pi_v| dy)^{\frac{5}{4}} d\tau ,$$

where $r \in (0, t^{\frac{1}{2}})$, and

$$\begin{aligned} & Q_r(t,x) := \{(\tau,y) : t - r^2 < \tau < t \text{ and } |y - x| < r\}, \\ & Q_r^*(t,x) := \{(\tau,y) : t - \frac{7}{8}r^2 < \tau < t + \frac{1}{8}r^2 \text{ and } |y - x| < r\}. \end{aligned}$$

In [CKN82] it is proved:

Proposition 1

Let (v, π_v) be a suitable weak solution in some parabolic cylinder $Q_r(t, x)$. There exist $\varepsilon_1 > 0$ and $c_0 > 0$ independent of (v, π_v) such that, if $M(r) \le \varepsilon_1$, then

$$|\mathbf{v}(\tau, \mathbf{y})| \le c_1^{\frac{1}{2}} r^{-1}, \ a. \ e. \ in(\tau, \mathbf{y}) \in Q_{\frac{r}{2}}(t, \mathbf{x}),$$
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where $c_1 := c_0 \varepsilon_1^{\frac{2}{3}}$. In particular, v is regular in $Q_{\frac{r}{2}}(t, x)$.

Proposition 2

There exists a constant $\varepsilon_3 > 0$ with the following property. If (v, π_v) is a suitable weak solution in some parabolic cylinder $Q_r^*(t, x)$ and

$$\limsup_{r\to 0} r^{-1} \iint_{Q_r^*} |\nabla v|^2 dy d\tau \le \varepsilon_3 \,,$$

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Also in the case of their applications, in particular Theorem C and Theorem D, and its Corollary, the results of regularity express on the geometry of regular points, but a behavior in a neighborhood of t = 0 does not seem an immediate consequence.

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Let (v, π_v) be a suitable weak solution corresponding to $v_o \in J^2(\mathbb{R}^3)$. There exist absolute constants ε_1 , C_1 and C_2 such that if

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In the above statement ε_1 is the same as in Proposition 1.

We remark that for all $v_{\circ} \in J^2(\mathbb{R}^3)$, by virtue of the Hardy-Littlewood-Sobolev theorem, we have that $\|v_{\circ}\|_{2,w(x)} < \infty$ almost everywhere in $x \in \mathbb{R}^3$.

Hence, a data in L^2 inherently defines functional (3)₁.

But in condition (4) there is more, that is the smallness of the norm $\|\|v_0\|_{2,w(x)}\|_{L^{\infty}(B(x_0,R_0))}$.

We can find several sufficient conditions on $v_o \in \mathcal{J}^2(\mathbb{R}^3)$ such that assumption (4) is verified. For example, if we consider the assumption made in [CKN82] (in the Corollary), that is $v_o \in W^{1,2}(\mathbb{R}^3 \setminus B_R)$, then $\|v_o\|_{2,w(x)}$ is a continuous function of x that we can make small outside a suitable ball of radius $R' \geq R$.

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Hence, a data in L^2 inherently defines functional (3)₁.

But in condition (4) there is more, that is the smallness of the norm $||| v_{\circ} ||_{2,w(x)} ||_{L^{\infty}(B(x_0,R_0))}$.

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A first result and comments: On the proof of the regularity result

The proof of Theorem 2 is based on Proposition 1.

Also in [CKN82] the regularity result is based on Proposition 1, but our approach to the regularity via the Proposition 1 is different. Indeed, we prove a weighted energy relation (the norm is $||v(t)||_{2,w(x)}$) which holds for t > 0, provided that (4)₁ holds:

$$\int_{\mathbb{R}^{3}} \frac{|v(t,y)|^{2}}{|x-y|} dy + c(\mathscr{E}_{0}) \int_{0}^{t} \int_{\mathbb{R}^{3}} \frac{|\nabla v(\tau,y)|^{2}}{|x-y|} dy d\tau \le c \int_{\mathbb{R}^{3}} \frac{|v_{\circ}(y)|^{2}}{|x-y|} dy,$$
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for all t > 0, a.e. in $B(x_0, R_0)$.

Thanks to estimate (6) we can apply Proposition 1 by [CKN82] to deduce regularity.

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Thanks to estimate (6) we can apply Proposition 1 by [CKN82] to deduce regularity.

The partial regularity has a local character in the sense that condition (4) can be not satisfied on $\mathbb{R}^3 \setminus B(x_0, R_0)$, and the regularity is ensured a.e. in $B(x_0, R_0)$.

Another feature that expresses the local character of the result is the following. The pointwise behavior of our solution v is given in a neighborhood of $(0, x_0)$, and, further it is as the one of the solutions to the Stokes problem. As far as we know, this property is new.

For this reason our result improves the previous one (that is $\nabla v_o \in L^2(\mathbb{R}^3 \setminus B_R)$) proved in [CKN82], in the sense that our thesis gives a bound and a behavior in neighborhood of t = 0:

$$|v(t,x)| \leq c(v_\circ)t^{-\frac{1}{2}}, \ (t,x) \in (0,\infty) imes \mathbb{R}^3 \setminus B_{R'}.$$

In particular, if $B(x_0, R_0) \equiv \mathbb{R}^3$, then we get a new sufficient condition for the existence of global smooth solutions.

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In particular, if $B(x_0, R_0) \equiv \mathbb{R}^3$, then we get a new sufficient condition for the existence of global smooth solutions.

Of course estimate (5) also gives a pointwise asymptotic behavior of the solution for large *t*. However such a behavior is not optimal under the assumption $v_o \in J^2(\mathbb{R}^3)$. Indeed for a suitable weak solution the pointwise asymptotic behavior is of the kind $O(\|v_o\|_2 t^{-\frac{3}{4}})$, according to the fact that a weak solution becomes smooth for $t > T_0$, where T_0 is connected with the L^2 -norm of v_o , and the behavior is stated considering the L^2 -norm of the initial data.

Leray proves that a weak solution becomes smooth, that is $\|v(t)\|_{\infty} < \infty$, for all $t > T_0$, where $T_0 \le c \|v_0\|_2^4$.

Based on Proposition 1 and Proposition 2, in [CKN82] it is proved that if $\|\nabla v_{\circ}\|_{L^{2}(\mathbb{R}^{3}-B_{R})} < \infty$, then there exists a $R' \ge R$ such that $\|v(t)\|_{L^{\infty}(\mathbb{R}^{3}-B_{R'})} < \infty$ for all t > 0.

Further, in the light of our Theorem 2 we can claim:

$$|v(t,x)| \leq c(v_\circ)t^{-rac{1}{2}}, ext{ for all } (t,x) \in (0,\infty) imes \mathbb{R}^3 - B_{R'}.$$

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$$|v(t,x)| \leq c(v_\circ)t^{-\frac{1}{2}}, \text{ for all } (t,x) \in (0,\infty) imes \mathbb{R}^3 - B_{R'}.$$

Roughly speaking, following Leray and the results in [CKN82], we say that the possible *turbulence* of a weak solution not only appears in a finite time ($t < T_0$), but also in a bounded region of the space ($B_{R'}$), whose \mathscr{P}^1 -measure is null (this last is true as soon as a suitable weak solution exists).

Actually, the smallness of the data for large |x|, although given by means of integrability conditions outside a ball, preserves the regularity of the weak solution (as in the case of *"small data"*).

The aim of our note goes in the direction of the last claims. We prove that, not only the possible *turbulence* does not perturb the regularity of a weak solution in a neighborhood of *infinity*, but, if an asymptotic spatial behavior of the initial data v_{\circ} is given, then the same behavior holds for a suitable weak solution for all t > 0.

As far as we know, such a property to date was ensured only for small data, [Knightly66], [F.C. and P.Maremonti 04].

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Theorem 3

Let $v_{\circ} \in J^{2}(\mathbb{R}^{3})$ and, for some $\alpha \in [1,3)$ and $R_{0} > 0$, let be $|v_{\circ}(x)| \leq V_{\circ}|x|^{-\alpha}$, for $|x| > R_{0}$. Let (v, π_{v}) be a suitable weak solution to the Navier-Stokes Cauchy problem. Then, there exists a constant $M \geq 1$ such that

$$|v(t,x)| \le c(v_{\circ})|x|^{-\alpha}, \text{ for all } (t,x) \in (0,\infty) \times \mathbb{R}^3 \setminus B_{MR_0}, \tag{7}$$

where M is independent of v_{\circ} . Moreover,

$$|\mathbf{v}(t,\mathbf{x})| \leq \mathbf{c}(\mathbf{v}_{\circ})t^{-\frac{lpha}{2}}, (t,\mathbf{x}) \in (\mathcal{T}_{0},\infty) imes \mathbb{R}^{3}.$$

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We work on a given suitable weak solution.

Hence we do not construct new solutions.

We are able to give a representation formula in $\mathbb{R}^3 \setminus B_{M_0R_0}$, where the existence of M_0 is proved in the first result.

Then, we get for all t>0 and almost a.e. in $x\in \mathbb{R}^3\setminus B_{M_0R_0}$:

$$\begin{aligned} v(t,x) &= \mathbb{H}[v_{\circ}](t,x) + \mathbb{T}[v \otimes v](t,x) \,, \\ \mathbb{H}[v_{0}](t-s,x) &:= \int_{\mathbb{R}^{3}} H(t-\tau,y)v_{\circ}(y)dy =: w(t,x) \,, \\ \mathbb{F}[v \otimes v](t,x) &:= \int_{0}^{t} \int_{\mathbb{R}^{3}} v(\tau,y) \cdot \nabla T_{i}(t-\tau,x-y) \cdot v(\tau,y)dyd\tau \,. \end{aligned}$$

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For the first term there are no problems on the spatial asymptotic behavior.

For the second term we consider

V = (V - W) + W =: U + W.

Hence,

 $\mathbb{T}[v,v](t,x) = \mathbb{T}[u,u](t,x) + \mathbb{T}[u \otimes w](t,x) + \mathbb{T}[w \otimes u](t,x) + \mathbb{T}[w \otimes w](t,x).$

It is not difficult to discuss the terms which contain the linear part w.

As far as it concerns the term $\mathbb{T}[u, u]$, that we call non-smooth term, we employ the following estimate

$$\|u(t)\|_{L^2(\mathbb{R}^3\setminus B_R)}^2 \leq c(v_\circ)R^{-1}$$

This behavior derives from a suitable discussion of the generalized energy inequality (it is not trivial). Of course, an iterative procedure is needed to perform the same order of decay of the initial data:

$|v(t,x)| \leq c(v_{\circ})|x|^{-\alpha}, t > 0, |x| > M_0 R_0.$

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